

# Symmetries of periodic and free boundary $q$ -Whittaker measures

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joint work with Michael Wheeler

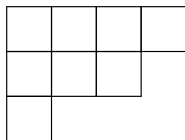
# Summary

This talk is about two things:

- Identities of symmetric functions ( $q$ -Whittaker and Hall–Littlewood functions)
- A different perspective on symmetric functions (connections to probability/statistical physics)

## Random partitions

# Schur measures



$$\lambda = (4, 3, 1), \quad l(\lambda) = 3.$$

- Many models of random partitions defined using [symmetric functions](#).
- Define the [Schur functions](#) by

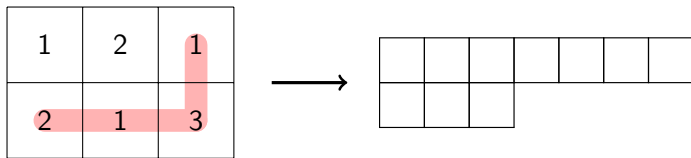
$$s_{\lambda}(x_1, \dots, x_n) = \frac{\det \left( x_i^{\lambda_j + n - j} \right)_{i,j=1}^n}{\prod_{i < j} (x_i - x_j)}.$$

This is a symmetric polynomial in the  $x_i$ 's.

- The [Schur measure](#) (Okounkov '01) with parameters  $x_1, \dots, x_n$  and  $y_1, \dots, y_m$  is probability measure on  $\lambda$  proportional to  $s_{\lambda}(x_1, \dots, x_n) s_{\lambda}(y_1, \dots, y_m)$ . Normalization constant given by [Cauchy identity](#)

$$\sum_{\lambda} s_{\lambda}(x_1, \dots, x_n) s_{\lambda}(y_1, \dots, y_m) = \prod_{i,j} \frac{1}{1 - x_i y_j}.$$

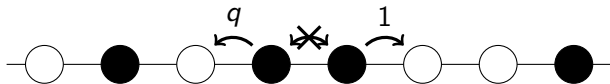
# Connections to statistical physics



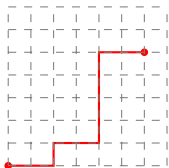
- For many choices of the  $x_i$  and  $y_j$ , the Schur measure is connected to models coming from statistical physics.
- The Robinson–Schensted–Knuth correspondence sends an  $n \times m$  array of integers to a pair of tableau of the same shape  $\lambda$ . If the integers are geometric random variables of parameter  $x_i y_j$ , then  $\lambda$  is distributed as the Schur measure.
- Define a model called **last passage percolation**. Consider a random field  $X_{i,j}$  of independent geometric random variables of parameter  $p$ . The **passage time** from  $(0,0)$  to  $(n,m)$  is given by  $\max_{\gamma} \sum_i X_{\gamma_i}$ , where  $\gamma$  is an up-right path. Under RSK this corresponds to  $\lambda_1$ .

# Macdonald measures

- Can define **Macdonald measure** using the **Macdonald polynomials**  $P_\lambda(x; q, t)$  analogously to the Schur measures (Borodin–Corwin '14). They have been used to study many models coming from statistical physics.
- When  $q = 0$ , this is called the Hall–Littlewood measure, related to the six vertex model and particle systems like the ASEP.



- When  $t = 0$ , this is called  $q$ -Whittaker measure, related to polymer models like the log-gamma polymer.



$$Z(6, 5) = \sum_{\gamma} \prod_i x_{\gamma_i}$$

## Periodic measures on partitions

# Periodic Schur measure

What is a periodic measure on partitions?

- Borodin '07 defined the **periodic Schur measure**, with the probability of  $\lambda$  proportional to  $\sum_{\mu} u^{|\mu|} s_{\lambda/\mu}(x_1, \dots, x_n) s_{\lambda/\mu}(y_1, \dots, y_m)$ . Here,  $\mu$  is a partition contained in  $\lambda$ , and  $s_{\lambda/\mu}$  is a **skew Schur function**.
- Normalization constant given by

$$\sum_{\lambda, \mu} u^{|\mu|} s_{\lambda/\mu}(x_1, \dots, x_n) s_{\lambda/\mu}(y_1, \dots, y_m) = \frac{1}{(u; u)_{\infty}} \prod_{i,j} \frac{1}{(x_i y_j; u)_{\infty}},$$

where  $(x; q)_{\infty} = \prod_{i \geq 0} (1 - q^i x)$  is the  **$q$ -Pochhammer symbol**.

- This measure turns out to be related to a **quasi-periodic** version of last passage percolation. One considers a cylinder tiled by  $n \times m$  arrays of geometric random variables, but where powers of  $u$  are accumulated in the parameters as well.



# Periodic $q$ -Whittaker measure

- We can define a probability measure on partitions called the **periodic  $q$ -Whittaker measure** by setting the probability of  $\lambda$  proportional to

$$\sum_{\mu} u^{|\mu|} P_{\lambda/\mu}(x_1, \dots, x_n; q, 0) Q_{\lambda/\mu}(y_1, \dots, y_n; q, 0).$$

The  $Q_{\lambda/\mu}$  are a scalar multiple of the  $P_{\lambda/\mu}$  (**skew Macdonald polynomial**).

- The normalization constant is

$$\frac{1}{(u; u)_{\infty}} \prod_{i,j} \frac{1}{(x_i y_j; u, q)_{\infty}},$$

where  $(x; q, u)_{\infty} = \prod_{i,j=0}^{\infty} (1 - x u^i q^j)$ . This is a version of the Cauchy identity.

# Symmetries of periodic $q$ -Whittaker measures

## Theorem (H.–Wheeler '23)

*The expression*

$$\sum_{\mu, \lambda: \lambda_1 \leq k} \frac{u^{|\mu|}}{(q; q)_{n-\lambda_1}} P_{\lambda/\mu}(x_1, \dots, x_n; q, 0) Q_{\lambda/\mu}(y_1, \dots, y_m; q, 0)$$

*is symmetric in the parameters  $u$  and  $q$ .*

- Generalizes an identity of Imamura–Mucciconi–Sasamoto '23, who showed the special case when one parameter is 0.
- The IMS proof is bijective, would be interesting to see if it could be extended to this case.

## Theorem (H.–Wheeler '23)

*The expression*

$$\sum_{\mu, \lambda: \lambda_1 \leq k} \frac{u^{|\mu|}}{(q; q)_{k-\lambda_1}} P_{\lambda/\mu}(x_1, \dots, x_n; q, 0) Q_{\lambda/\mu}(y_1, \dots, y_m; q, 0)$$

*is symmetric in the parameters  $u$  and  $q$ .*

- The expression has a probabilistic interpretation: If  $\chi$  is a  $q$ -geometric random variable, meaning  $\mathbb{P}(\chi \leq n) = \frac{(q; q)_\infty}{(q; q)_n}$ , then the expression is equivalent to  $\mathbb{P}(\lambda_1 + \chi \leq k)$ .

# Contour integral formulas

## Theorem (H.–Wheeler '23)

*The expression*

$$\sum_{\mu, \lambda: \lambda_1 \leq k} \frac{u^{|\mu|}}{(q; q)_{n-\lambda_1}} P_{\lambda/\mu}(x_1, \dots, x_n; q, 0) Q_{\lambda/\mu}(y_1, \dots, y_m; q, 0)$$

*equals*

$$\Phi_n(u, q) \oint_C \frac{dz_1}{2\pi iz_1} \cdots \oint_C \frac{dz_k}{2\pi iz_k} \prod_{i,j} (1 + z_i^{-1} x_j) \prod_{i,j} (1 + z_i y_j) \tilde{\Delta}(z; q, u),$$

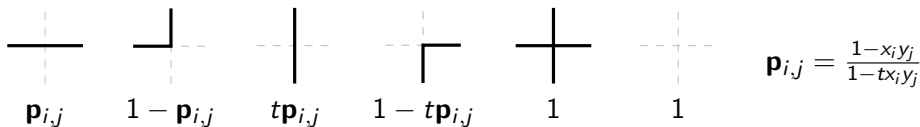
*where  $C$  is a circle centered at 0,*

$$\Phi_n(u, q) = \frac{(1 - qu)^k}{n!(1 - u)^k(1 - q)^k},$$

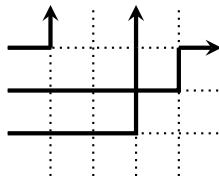
*and*

$$\tilde{\Delta}(z; q, u) = \prod_{i \neq j} \frac{(1 - quz_i z_j^{-1})(1 - z_i z_j^{-1})}{(1 - qz_i z_j^{-1})(1 - uz_i z_j^{-1})}.$$

# Stochastic six vertex model

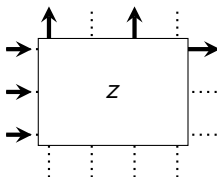


- Stochastic six vertex model is probability distribution on configurations of arrows traveling edges of a grid.
- Model depends on global parameter  $t$ , and row/column parameters  $x_1, \dots, x_n$  and  $y_1, \dots, y_m$ .
- Can be sampled one vertex at a time. Probabilities at vertex  $(i, j)$  depend on  $t$  and  $x_i y_j$  (called the **spectral parameter**). Come from  $R$ -matrix of  $U_q(\widehat{\mathfrak{sl}}_2)$ .

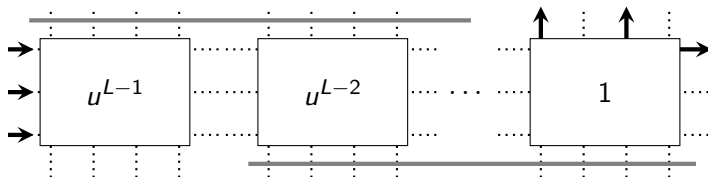


# Quasi-periodic boundary conditions

We represent a single  $m \times n$  6VM (with spectral parameter scaled by  $z$ ) by



We define **quasi-periodic 6VM** of length  $L$  by



Gray line indicates identification of edges. Can send  $L \rightarrow \infty$ , and we do so.

# Periodic Hall–Littlewood measure

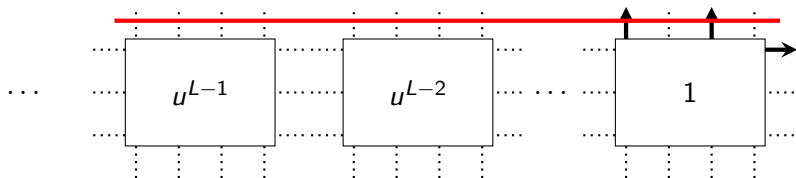
- Analogous to periodic  $q$ -Whittaker measure, we define the **periodic Hall–Littlewood measure** by setting probability of  $\lambda$  proportional to

$$\sum_{\mu} u^{|\mu|} P_{\lambda/\mu}(x_1, \dots, x_n; 0, t) Q_{\lambda/\mu}(y_1, \dots, y_m; 0, t).$$

- Normalization constant is

$$\frac{1}{(u; u)_{\infty}} \prod_{i,j} \frac{(tx_i y_j; u)_{\infty}}{(x_i y_j; u)_{\infty}}.$$

# Quasi-periodic 6VM and periodic HL measure



## Theorem (H.–Wheeler '23)

*The distribution for the number of times **arrows cross the red line** plus  $\chi$  equals that of  $I(\lambda)$ , where  $\lambda$  follows the periodic HL measure and  $\chi$  is  $u$ -geometric (with parameters in both models matching).*

- Proof uses vertex models, Yang–Baxter equation.
- When  $u \rightarrow 1$ , model becomes stationary **periodic** stochastic six vertex model, but observable in theorem blows up.



Free boundary measures

# Free boundary measures

- So far, we've discussed various measures related to the **Cauchy identity**. We now turn to the **Littlewood identity**, which turns out to correspond to models with either one or two open boundaries.
- The classic example is the **Pfaffian Schur measure**, for which  $\lambda$  has probability proportional to  $s_\lambda(x_1, \dots, x_n) \mathbf{1}_{\lambda' \text{ even}}$  (here  $\lambda'$  means rows and columns in picture are swapped). Normalization constant given by the Littlewood identity

$$\sum_{\lambda: \lambda' \text{ even}} s_\lambda(x_1, \dots, x_n) = \prod_{i < j} \frac{1}{1 - x_i x_j}.$$

- Generalized by Betea–Bouttier–Nejjar–Vuletić '18 to **free boundary Schur measure**, where probability of  $\lambda$  proportional to  $\mathbf{1}_{\lambda' \text{ even}} \sum_{\mu' \text{ even}} u^{|\mu|} s_{\lambda/\mu}(x_1, \dots, x_n)$ , with normalization constant

$$\frac{1}{(u; u)_\infty} \prod_{i < j} \frac{1}{(x_i x_j; u)_\infty}$$

# Free boundary $q$ -Whittaker measure

- We can define a probability measure on partitions called the **periodic  $q$ -Whittaker measure** by setting the probability of  $\lambda$  proportional to

$$\sum_{\mu} u^{|\mu|/2} h_{\lambda'}^*(a, b; q) h_{\mu'}(c/\sqrt{u}, d/\sqrt{u}; q) P_{\lambda/\mu}(x_1, \dots, x_n; q, 0).$$

The  $h_{\lambda}^*$  and  $h_{\mu}$  are **boundary weights**, defined in terms of parameters  $a, b, c, d$ . If  $a = b = c = d = 0$ , restrict  $\mu'$  and  $\lambda'$  to be even.

- The normalization constant is explicit and factorized (but complicated). When  $a = b = c = d = 0$ , it's

$$\frac{1}{(u; u)_{\infty} (uq; u, q)_{\infty}} \prod_{i < j} \frac{1}{(x_i x_j; u, q)_{\infty}}.$$

# Symmetries of free boundary $q$ -Whittaker measures

## Theorem (H.–Wheeler '25+)

*The expression*

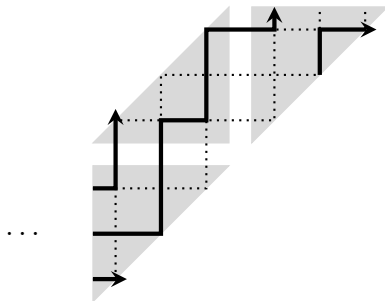
$$\sum_{\lambda, \mu: \lambda_1 \leq k} u^{|\mu|/2} \frac{h_{k-\lambda_1}(ab; q) h_{\lambda'}^*(a, b; q)}{(q; q)_{k-\lambda_1}} h_{\mu'}(c/\sqrt{u}, d/\sqrt{u}; q) P_{\lambda/\mu}(x; q, 0)$$

*is symmetric in the parameters  $u$  and  $q$ , and separately in  $a, b, c, d$ .*

- As in periodic setting, we find explicit contour integral formulas, and this has a probabilistic interpretation.
- Generalizes previous identities of Imamura–Mucciconi–Sasamoto '23 (and extended in H. '24).

# Free boundary Hall–Littlewood measure

- We can define the **free boundary Hall–Littlewood measure** analogously. Like in periodic setting, can be related to observables in **quasi-open six vertex model**. Here arrows enter from diagonal with weights depending on  $a, b, c, d$ .



# A pair of curious identities

## Theorem (H.–Wheeler '23)

$$\sum_{\lambda, \mu: \lambda_1 \leq k} \frac{(t; t)_{m_k(\mu)}}{(t; t)_{m_k(\lambda)}} u^{|\mu|} P_{\lambda/\mu}(x_1, \dots, x_n; 0, t) Q_{\lambda/\mu}(y_1, \dots, y_m; 0, t) \\ = \frac{1}{(u; u)_k} \left( \prod_{j=1}^m y_j^k \right) P_{k^m}(x_1, \dots, x_n, y_1^{-1}, \dots, y_m^{-1}; u, t).$$

## Theorem (Finn–Vanicat '17)

$$\sum_{\lambda, \mu: \lambda_1 \leq 2k} u^{|\mu|/2} \frac{h_\lambda(a, b; t)}{h_{m_{2k}(\lambda)}(ab; t)} h_\mu^*(c/\sqrt{u}, d/\sqrt{u}; t) P_{\lambda/\mu}(x_1, \dots, x_n; 0, t) \\ = C_k(u, t, a, b, c, d) \left( \prod_{i=1}^n x_i^k \right) K_{k^n}(x_1, \dots, x_n; u, t, a, b, c, d).$$

Thanks for listening!