#### Lusztig's q-weight multiplicities and KR crystals

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(joint work with Donghyun Kim and Seung Jin Lee) arxiv.org/abs/2412.20757

## Symmetric function

For a partition  $\lambda$ , the Schur function  $s_{\lambda}$  is defined by

$$s_{\lambda} = \sum_{T} x^{\mathsf{wt}(T)}$$

where the sum is over all semistandard Young tableaux T of shape  $\lambda$ , and for a composition  $\alpha=(\alpha_1,\alpha_2,\dots)$ , we write  $x^\alpha=x_1^{\alpha_1}x_2^{\alpha_2}\cdots$ .

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The Schur function can be expressed as  $s_{\lambda} = \sum_{\mu} K_{\lambda,\mu} m_{\mu}$ , where  $m_{\mu}$  is the monomial symmetric function and  $K_{\lambda,\mu}$  is the Kostka number.

The Kostka number  $K_{\lambda,\mu}$  counts the number of semistandard Young tableaux of shape  $\lambda$  and weight  $\mu$ .

## Kostka-Foulkes polynomial

The Kostka-Foulkes polynomials  $K_{\lambda,\mu}(q)$ , q-analogue of the Kostka number, are defined by

$$s_{\lambda} = \sum_{\mu} K_{\lambda,\mu}(q) P_{\mu}(x;q),$$

where  $P_{\mu}(x;q)$  is the Hall-Littlewood polynomial.

The charge is a statistic on semistandard Young tableaux, introduced by Lascoux and Schtzenberger (1978). It provides a combinatorial formula for the Kostka-Foulkes polynomial:

$$K_{\lambda,\mu}(q) = \sum_{T} q^{\operatorname{charge}(T)}$$

where the sum is over all semistandard Young tableaux T of shape  $\lambda$  and weight  $\mu.$ 

We define the charge on a standard Young tableau T.

Let  $\mathrm{Des}(T)$  be the set of integers i such that i+1 appears to the right of i in T. Then the charge is defined by  $\sum_{i\in\mathrm{Des}(T)}(n-i)$ .

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For example, consider the standard Young tableau

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For example, consider the standard Young tableau

$$T = \begin{bmatrix} 1 & 3 & 6 \\ 2 & 4 \\ 5 \end{bmatrix}.$$

Then  $Des(T) = \{2, 5\}$  and charge(T) = (6-2) + (6-5) = 5.

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- Lusztig's q-weight multiplicities, which generalize the Kostka-Foulkes polynomials to other Lie types.

## Weight multiplicity

For a Lie algebra  $\mathfrak{g}$ , let  $\lambda$  and  $\mu$  be dominant weights.

The weight multiplicity  $KL_{\lambda,\mu}^{\mathfrak{g}}$  is the dimension of the  $\mu$ -weight space in the irreducible representation of  $\mathfrak{g}$  with highest weight  $\lambda$ .

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By the Weyl character formula,

$$\mathrm{KL}_{\lambda,\mu}^{\mathfrak{g}} = \sum_{w \in W} (-1)^{\ell(w)} [e^{w(\lambda+\rho)-(\mu+\rho)}] \prod_{\alpha \in R^+} \frac{1}{1-e^{\alpha}}$$

where W is Weyl group,  $R^+$  is the set of positive roots,  $\rho=\frac{1}{2}\sum_{\alpha\in R^+}\alpha$ , and  $[e^\beta]f$  denotes the coefficient of  $e^\beta$  in f.

The q-analogue of  $\mathrm{KL}_{\lambda,\mu}^{\mathfrak{g}}$  is defined by

$$\mathrm{KL}_{\lambda,\mu}^{\mathfrak{g}}(q) = \sum_{w \in W} (-1)^{\ell(w)} \left[ e^{w(\lambda + \rho) - (\mu + \rho)} \right] \prod_{\alpha \in \mathbb{R}^+} \frac{1}{1 - qe^{\alpha}}.$$

 $\mathrm{KL}_{\lambda,\mu}^{\mathfrak{g}}(q)$  is Lusztig's q-weight multiplicity.

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- $\mathrm{KL}_{\lambda,\mu}^{\mathfrak{g}}(q) = q^{(\ell(\omega_{\lambda}) \ell(\omega_{\mu})/2)} P_{\omega_{\mu},\omega_{\lambda}}^{\hat{\mathfrak{g}}}(q^{-1})$ , where  $P_{x,y}^{\hat{\mathfrak{g}}}(q)$  is affine Kazhdan-Lusztig polynoimal. So,  $\mathrm{KL}_{\lambda,\mu}^{\mathfrak{g}}(q) \in \mathbb{Z}_{\geq 0}[q]$ . [Lusztig '83].

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- $\mathrm{KL}^{\mathfrak{g}}_{\lambda+(k^n),\mu+(k^n)}(q)$  stabilize for sufficiently large k. These are called stable  $\mathrm{KL}$  polynomials.

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- Combinatorial formulas for stable KL polynomials are known: for types B and C by Shimozono (2005); and for type D by Lecouvey and Shimozono (2007).

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- Lecouvey (2005) studied cyclage graphs on Kashiwara-Nakashima tableaux and defined charge to conjecture a combinatorial formula.
- Lecouvey and Lenart (2020) provided a combinatorial formula when weight is zero, i.e.  $\mathrm{KL}_{\lambda,0}^{C_n}(q)$ .

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For a partition  $\mu$ , Lee conjectured a formula in terms of Killirov-Reshitikhin crystals, which was proved by C.-Kim-Lee (2024).

## Combinatorial object

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- Kashiwara-Nakashima tableaux.
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- Semistandard oscillating tableaux [Lee '23].

In this talk, we focus on semistandard oscillating tableaux.

An oscillating horizontal strip (ohs) S is a triple of partitions  $(\lambda,\mu,\nu)$  such that both  $\mu/\lambda$  and  $\mu/\nu$  are horizontal strips.

#### We define:

- length(S) =  $|\mu/\lambda| + |\mu/\nu|$ ,
- $I(S) = \lambda$ ,  $F(S) = \nu$ , and  $c(S) = \mu_1$ .

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A semistandard oscillating tableaux T is a sequence of ohs

$$(S_1, S_2, \cdots, S_n)$$
 satisfying:

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- $I(S_{i+1}) = F(S_i)$  for  $i \in \mathbb{Z}_{>0}$ .

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We define the following:

- $\operatorname{wt}(T) = (\operatorname{length}(S_1), \dots, \operatorname{length}(S_n))$  and  $c(T) = \max(c(S_i))$ .
- $T \in SSOT(\lambda, \mu)$  if  $F(S_n) = \lambda$  and  $wt(T) = \mu$ .
- $T \in SSOT_{\leq g}(\lambda, \mu)$  if  $T \in SSOT(\lambda, \mu)$  and  $c(T) \leq g$ .

## Example

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 $SSOT(\lambda, \mu)$  consists of the following:

- $T_1 = ((\emptyset, \square, \square), (\square, \square, \square), (\square, \square, \square), (\square, \square, \emptyset))$
- $T_2 = ((\emptyset, \square, \square), (\square, \square, \emptyset), (\emptyset, \square, \square), (\square, \square, \emptyset)),$
- $T_3 = ((\emptyset, \square, \square), (\square, \square, \square), (\square, \square, \square), (\square, \square, \emptyset)).$

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$$T_3 = ((\emptyset, \square, \square), (\square, \square), (\square, \square), (\square, \square, \square), (\square, \square, \emptyset)).$$

 $\mathrm{SSOT}_{\leq 1}(\lambda,\mu)$  contains only  $T_1$  and  $T_2$ , since  $c(T_3)=2$ .

## Type C object

Lee (2023) proved that for any  $g \ge \lambda_1$ ,

$$\mathrm{KL}_{\lambda,\mu}^{C_n} = |\operatorname{SSOT}_{\leq g}(\hat{\lambda}, \hat{\mu})|,$$

where 
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We now investigate the energy function on the Kirillov–Reshetikhin crystal, which plays the role of the charge statistic on SSOT.

#### Kirillov-Reshetikhin crystal

The Kirillov-Reshetikhin crystals (KR crystals) are crystal bases  $B^{r,s}$  for certain irreducible finite-dimensional modules  $W_s^{(r)}$ , called Kirillov-Reshetikhin modules (KR modules), over the quantized affine algebra  $U_a'(\mathfrak{g})$ .

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- For a partition  $\mu$ , define  $B_{\mu} = B^{\mu_n,1} \otimes \cdots \otimes B^{\mu_1,1}$ .
- For a partition  $\lambda$ , let  $HW(B_{\mu}, \lambda)$  be the set of classical highest weight elements of weight  $\lambda^t$  in  $B_{\mu}$ .
- $\bullet$  By duality, we use KR crystals of type  $B_N^{(1)}$  for sufficiently large N.

#### Kirillov-Reshetikhin crystal

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For  $SSOT_{<1}((0,0,0,0),(1,1,1,1))$ , we have:

$$T_{1} = ((\emptyset, \square, \square), (\square, \square, \square), (\square, \square, \square), (\square, \square, \emptyset)) \leftrightarrow -1 \otimes -1 \otimes 1 \otimes 1$$

$$T_{2} = ((\emptyset, \square, \square), (\square, \square, \emptyset), (\emptyset, \square, \square), (\square, \square, \emptyset)) \leftrightarrow -1 \otimes 1 \otimes -1 \otimes 1$$

$$T_2 = ((\emptyset, \square, \square), (\square, \square, \emptyset), (\emptyset, \square, \square), (\square, \square, \emptyset)) \leftrightarrow -1 \otimes 1 \otimes -1 \otimes 1$$

## **Energy function**

The energy function is a map  $\overline{D}:B\to\mathbb{Z}$ , which is constant on each classical component.

It is defined locally using the combinatorial R-matrix and the crystal operator  $e_0$ , and then extended globally.

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It is defined locally using the combinatorial R-matrix and the crystal operator  $e_0$ , and then extended globally.

In general, the energy function is difficult to compute.

However, when  $B=(B^{1,1})^{\otimes n}$ , the energy function can be computed explicitly.

## Energy function when standard case

Let  $a_n \otimes a_{n-1} \otimes \cdots \otimes a_1$  be an element in  $(B^{1,1})^{\otimes n}$  of type  $B_N^{(1)}$ .

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$$H(b,a) = \begin{cases} 2 & \text{if } a = 1 \text{ and } b = \overline{1} \\ 1 & \text{if } b \succ a \text{ and } (b,a) \neq (\overline{1},1) \\ 0 & \text{if } b \preceq a \end{cases}$$

under the order  $1 \prec 2 \prec \cdots \prec \bar{2} \prec \bar{1}$ .

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under the order  $1 \prec 2 \prec \cdots \prec \bar{2} \prec \bar{1}$ . The energy function  $\overline{D}$  is defined by  $\overline{D} = \sum\limits_{i=1}^{n-1} (n-i)H(a_{i+1},a_i)$ .

### Type A case

For type A, the energy function coincides with the charge statistics [Nakayashiki-Yamada '97].

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For example, consider

$$T = \begin{array}{|c|c|}\hline 1 & 3 & 6 \\\hline 2 & 4 \\\hline 5 & \end{array} \leftrightarrow 3 \otimes 1 \otimes 2 \otimes 2 \otimes 1 \otimes 1$$

Since  $3\otimes 1\otimes 2\otimes 2\otimes 1\otimes 1$ , we have  $\overline{D}(T)=(6-5)+(6-2)=5$ , which exactly matches the charge of T.

SSOT(
$$(0,0,0,0),(1,1,1,1)$$
) =  $\{T_1,T_2,T_3\}$  where  $T_1 = -1 \otimes -1 \otimes 1 \otimes 1$ ,  $T_2 = -1 \otimes 1 \otimes -1 \otimes 1$ , and  $T_3 = -1 \otimes -2 \otimes 2 \otimes 1$ .

$$\begin{split} & \text{SSOT}((0,0,0,0),(1,1,1,1)) = \{T_1,T_2,T_3\} \text{ where } T_1 = -1 \otimes -1 \otimes 1 \otimes 1, \\ & T_2 = -1 \otimes 1 \otimes -1 \otimes 1, \text{ and } T_3 = -1 \otimes -2 \otimes 2 \otimes 1. \end{split}$$

- $-1 \otimes -1 \otimes 1 \otimes 1$  with  $\overline{D}(T_1) = (4-2) \times 2 = 4$ .
- $-1 \otimes 1 \otimes -1 \otimes 1$  with  $\overline{D}(T_2) = (4-3) \times 2 + (4-1) \times 2 = 8$ .
- $-1 \otimes -2 \otimes 2 \otimes 1$  with  $\overline{D}(T_3) = (4-3) + (4-2) + (4-1) = 6$ .

#### Main theorem

#### Theorem (C.-Kim-Lee, 2024)

$$\mathrm{KL}_{\lambda,\mu}^{C_n}(q) = \sum_{T \in \mathrm{SSOT}_{\leq q}(\hat{\lambda},\hat{\mu})} q^{\overline{D}(T)},$$

where  $g \geq \lambda_1$ .

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where  $q > \lambda_1$ .

As a corollary, we have  $\mathrm{KL}_{\lambda+(1^n),\mu+(1^n)}^{C_n}(q) \geq \mathrm{KL}_{\lambda,\mu}^{C_n}(q)$ ,

since  $SSOT_{\leq g+1}(\hat{\lambda}, \hat{\mu}) \supseteq SSOT_{\leq g}(\hat{\lambda}, \hat{\mu})$ .

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When g=1, we have  $\lambda=(1,1,1,1)$  and  $\mu=(0,0,0,0).$ 

$$\mathrm{KL}_{(1,1,1,1),(0,0,0,0)}^{C_n}(q) = q^{\overline{D}(T_1)} + q^{\overline{D}(T_2)} = q^8 + q^4.$$

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When g=2, we have  $\lambda=(2,2,2,2)$  and  $\mu=(1,1,1,1)$ .

$$\mathrm{KL}_{(2,2,2,2),(1,1,1,1)}^{C_n}(q) = q^{\overline{D}(T_1)} + q^{\overline{D}(T_2)} + q^{\overline{D}(T_3)} = q^8 + q^6 + q^4.$$

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The level-restricted q-weight multiplicity  $\mathrm{KL}_{\lambda,\mu}^{\mathfrak{g},lr}(q)$  is defined by

$$\mathrm{KL}_{\lambda,\mu}^{\mathfrak{g},lr}(q) = \sum_{w \in W} (-1)^{\ell(w)} [e^{w(\lambda+\rho)-(\mu+\rho)}] \prod_{\alpha \in R_A} \frac{1}{1-qe^{\alpha}} \prod_{\alpha \in R^+ \backslash R_A} \frac{1}{1-e^{\alpha}}.$$

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We also proved the following formula

$$\mathrm{KL}_{\lambda,\mu}^{C_n,lr}(q) = \sum_{\substack{T \in \mathrm{SSOT}(\hat{\lambda},\hat{\mu}) \\ c(T) \leq g}} q^{||\hat{\mu}|| + \frac{|\hat{\mu}| - |\hat{\lambda}|}{2} - \overline{D}(\phi_r(T))}$$

using the row KR crystals  $B^{1,\mu_n} \otimes \cdots \otimes B^{1,\mu_1}$  of type  $C_N^{(1)}$ .

## Summary and future direction

We also investigate these multiplicities for other Lie types.

	Lusztig's $q$ -weight multiplicity		I.r. q-weight multiplicity
type A	Lascoux and Schützenberger (1978)		
type B	$D_{N+1}^{(2)}$ -column	?	$D_{N+1}^{(2)}$ -row
type C	$B_N^{(1)}$ -co	olumn	$C_N^{(1)}$ -row
type D	?		$B_N^{(1)}$ -row

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- A natural next step is to fill in the missing entries.
- It would also be interesting to investigate the connection with rigged configurations.

# Thanks for listening