

THE PENNEY-FUJIWARA PLANCHEREL FORMULA FOR HOMOGENEOUS SPACES WITH MONOMIAL SPECTRUM

RONALD L. LIPSMAN

Department of Mathematics, University of Maryland

The Penney-Fujiwara Plancherel Formula [PFPF] is a distribution-theoretic version of the direct integral decomposition of the quasi-regular representation $\tau = \text{Ind}_H^G 1$. Here G is a connected Lie group and H is a closed subgroup. Aside from its intrinsic interest, such a formula has the advantage over an abstract representation-theoretic decomposition in that it yields the explicit intertwining operator that effects the decomposition. The main feature of [3] and this short note is the presentation of such a formula in several general situations in which the generic representations that occur in the decomposition are monomial.

The main results of [3] can be summarized as follows. We fix right Haar measures dg, dh on G, H and a smooth function $q = q_{H,G}$ on G satisfying $q_{H,G}(hg) = \frac{\Delta_H(h)}{\Delta_G(h)} q_{H,G}(g)$, where Δ denotes modular functions. These choices uniquely determine a quasi-invariant measure $d\dot{g}$ on $H \backslash G$ satisfying

$$\int_G f(g)q(g)dg = \int_{H \backslash G} \int_H f(hg)dhd\dot{g}$$

Suppose it is true generically that, for each irreducible representation π that occurs in τ , there is a closed subgroup B of G and a character χ of B for which $\pi = \pi_\chi = \text{Ind}_B^G \chi$. We realize the representation π_χ in the space $\mathcal{H}_\chi = \{f : G \rightarrow \mathbb{C}, f(bg) = \chi(b)f(g), \int_{B \backslash G} |f|^2 d\dot{g} < \infty\}$, where $d\dot{g}$ is the quasi-invariant measure on $B \backslash G$ now, and G acts by $\pi_\chi(g)f(x) = f(xg)[q_{B,G}(xg)/q_{B,G}(x)]^{1/2}$. We further assume that BH is closed and $q_{H \cap B, H} q_{H \cap B, B} \equiv 1$ on $H \cap B$. Under these conditions we can define a (conjugate-linear) distribution

$$\beta_\chi : f \rightarrow \int_{H \cap B \backslash H} \bar{f} q_{B,G}^{1/2} q_{H \cap B, H}^{-1} q_{H,G}^{-1/2} dh, f \in C_c^\infty(G, B, \chi).$$

Then the following facts are true.

- (i) β_χ is relatively invariant under the action of H with modulus $q_{H,G}^{-1/2}$.
- (ii) $\pi_\chi(\omega)\beta_\chi \in C^\infty(G, B, \chi)$ for any test function ω on G ; in fact

$$\pi_\chi(\omega)\beta_\chi(g) = \int_{H \cap B \backslash B} \omega_H(bg) \bar{\chi}(b) q_{B,G}^{-1/2}(bg) q_{H,G}^{1/2}(bg) q_{H \cap B, B}^{-1}(b) db,$$

where

$$\omega_H(g) = \Delta_G(g)^{-1} q_{H,G}^{-1/2}(g) \int_H \omega(g^{-1}h^{-1}) \Delta_G(h)^{-1} q_{H,G}^{-1/2}(h) dh.$$

(iii) The matrix coefficient $\langle \pi_\chi(\omega) \beta_\chi, \beta_\chi \rangle$ of β_χ is well-defined and equals

$$(1) \int_{H \cap B \setminus H} \int_{H \cap B \setminus B} \omega_H(bh) \bar{\chi}(b) q_{B,G}^{-1/2}(b) q_{H,G}^{1/2}(h^{-1}bh) q_{H \cap B, B}^{-1}(b) q_{H \cap B, H}^{-1}(h) db dh.$$

I wish to emphasize that these results require NO structural assumptions on G other than those already stipulated. Now suppose we have a nice parameter space for the characters χ , say Λ . The PFPF amounts to the elucidation of a measure $d\chi$ on Λ for which one can prove the integral formula

$$(2) \int_{\Lambda} \langle \pi_\chi(\omega) \beta_\chi, \beta_\chi \rangle d\chi = \langle \tau(\omega) \alpha_\tau, \alpha_\tau \rangle.$$

Here α_τ is the canonical (relatively invariant) distribution on $\mathcal{H}_\tau^\infty$, and the right side of (2) equals $\omega_H(e)$ (see [3]). It is possible to deduce from formula (2) (see [3]) that the intertwining operator for the direct integral decomposition

$$\tau = \int_{\Lambda}^{\oplus} \pi_\chi d\chi$$

is

$$\Omega(g) \rightarrow \int_{H \cap B \setminus B} \Omega(bg) \bar{\chi}(b) q_{B,G}^{-1/2}(bg) q_{H,G}^{1/2}(bg) q_{H \cap B, B}^{-1}(b) db$$

where Ω is any H -invariant test function on G .

When can we actually carry out this program? That is, for which homogeneous spaces G/H whose spectrum is monomial can we actually progress from the computation of the matrix coefficients (1) to the PFPF (2)? [3] contains two general situations in which this is accomplished: (a) G simply connected completely solvable, H connected; and (b) G/H an abelian symmetric space, i.e. G is a semidirect product $G = HV$ where V is a closed normal vector subgroup and H is arbitrary (other than $\hat{V}/H$ is countably separated). In this note I would like to announce several other general situations where we can derive the PFPF from the matrix coefficients. I will list three.

(1) G exponential solvable. The work here generalizes (a) above precisely as Fujiwara's theory in [1] subsumes that of [2].

(2) $G = RV$, where R is connected reductive, V is a normal vector subgroup and $H = SU$ is of the same form with $S \subset R, U \subset V$. Furthermore, suppose the generic stabilizer in R of elements from $U^\perp$ is trivial. In this case the parameter space $\Lambda = U^\perp/S$, and there is a common polarization for all $\chi \in \Lambda$, namely $B = V$. Moreover, there is a natural realization of the Plancherel measure as a pseudo-image of Lebesgue measure on $U^\perp$, depending on the choices of the q functions.

(3) Start with $G = SL(2, \mathbb{R}) \cdot \mathbb{R}^2$. G may be realized as the group of real matrices of the form

$$G = \begin{pmatrix} a & b & x \\ c & d & y \\ 0 & 0 & 1 \end{pmatrix}, ad - bc = 1.$$

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Conjugation of G by the diagonal matrix with entries $(1, -1, 1)$ is an involution. The stability subgroup is

$$H = \begin{pmatrix} a & 0 & x \\ 0 & a^{-1} & 0 \\ 0 & 0 & 1 \end{pmatrix}.$$

G/H is a symmetric space and $\tau = \int^{\oplus} \pi_t dt$ where $\pi_t = \text{Ind}_B^G \chi_t$,

$$B = \begin{pmatrix} 1 & b & x \\ 0 & 1 & y \\ 0 & 0 & 1 \end{pmatrix},$$

and $\chi_t(b, x, y) = e^{ibt} e^{iy}$. In this case Lebesgue measure dt is the Plancherel measure and the derivation of the PFPF involves a Fourier inversion argument.

REMARKS. (i) Number (3) is representative of a class of semidirect products (of semisimple by vector groups) for which I can derive the PFPF. The key feature of those (as well as the spaces in number 2) is not only that the generic representations are monomial, but that there is a common polarization which is normalized by H .

(ii) A very interesting situation in which one has a common real polarization not normalized by H is that of a Riemannian symmetric space G/K (G a semisimple Lie group, K a maximal compact subgroup). The common polarization is a minimal parabolic subgroup. It is possible that one might be able to simplify the usual presentation of the spherical Plancherel measure by the techniques described here.

(iii) I have also considered situations in which there exist common positive (non-real) totally complex polarizations—especially various oscillator groups. I will discuss these at the conference as well as in another manuscript.

REFERENCES

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